

1. a) does wavelength compare favorably with atomic spacing in crystals ($\sim 10^{-10}$ m)?

$$\lambda = \frac{h}{p} = \frac{h}{mv} \quad \text{thermal speed } v \ll c, \text{ relativity not needed}$$

$$\text{thermal speed: } \text{cmv. } E \quad K = \frac{1}{2}mv^2 = E_{\text{th}} = \frac{3}{2}k_B T$$

$$\Rightarrow v = \sqrt{\frac{3k_B T}{m}} \approx 2.7 \times 10^3 \text{ m/s} \quad \text{classical expression justified}$$

$v \ll c$

$$\Rightarrow \lambda = \frac{h}{mv} = \frac{h}{m} \sqrt{\frac{m}{3k_B T}} = \frac{h}{\sqrt{3k_B T m}} \approx 1.5 \times 10^{-10} \text{ m}$$

$\lambda \sim$ same order as crystal spacing, it will work

b) $E = 1 \text{ GeV} \gg mc^2$ for electron, will need relativity

$$K = \sqrt{p^2 c^2 + m^2 c^4} - mc^2$$

$$\Rightarrow p^2 = \frac{(K + mc^2)^2 - m^2 c^4}{c^2}$$

$$\lambda = \frac{h}{p} \approx 1.2 \times 10^{-15} \text{ m}$$

of order H atom nucleus diameter

2. need $\frac{\partial \psi}{\partial x}$, ψ^* , $\frac{\partial \psi^*}{\partial x}$. presume A, B could be complex ...

$$\frac{\partial \psi}{\partial x} = \frac{i p}{\hbar} \left[A e^{i p x / \hbar} - B e^{-i p x / \hbar} \right] e^{-i p^2 t / 2 m \hbar}$$

$$\psi^* = \left[A^* e^{-i p x / \hbar} + B^* e^{i p x / \hbar} \right] e^{i p^2 t / 2 m \hbar}$$

$$\frac{\partial \psi^*}{\partial x} = \frac{i p}{\hbar} \left[-A^* e^{-i p x / \hbar} + B^* e^{i p x / \hbar} \right] e^{i p^2 t / 2 m \hbar}$$

$$\begin{aligned} \psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} &= \frac{i p}{\hbar} \left[A A^* - A^* B e^{-2 i p x / \hbar} + A B^* e^{2 i p x / \hbar} - B B^* \right] \\ &\quad - \frac{i p}{\hbar} \left[-A A^* - A^* B e^{-2 i p x / \hbar} + A B^* e^{2 i p x / \hbar} + B B^* \right] \end{aligned}$$

$$\psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} = \frac{i p}{\hbar} [2 A^* A - 2 B^* B] = \frac{2 i p}{\hbar} [|A|^2 - |B|^2]$$

$$j = \frac{\hbar}{2 m i} \cdot \frac{2 i p}{\hbar} [|A|^2 - |B|^2] = \frac{p}{m} [|A|^2 - |B|^2] = v (|A|^2 - |B|^2)$$

with $p = m v$

[note $\rho = |\psi|^2 = |A|^2 - |B|^2$, and we have $j = \rho v$
like an ideal incompressible fluid]

$$3. \quad \frac{\partial \rho}{\partial t} = \frac{\partial}{\partial t} (\psi^* \psi) = \frac{\partial \psi^*}{\partial t} \psi + \psi^* \frac{\partial \psi}{\partial t}$$

$$\frac{\partial \psi}{\partial t} = -\frac{\hbar}{2mi} \frac{\partial^2 \psi}{\partial x^2} + \frac{1}{i\hbar} V \psi \quad \text{time-dep Schrödinger}$$

$$\frac{\partial \psi^*}{\partial t} = \frac{\hbar}{2mi} \frac{\partial^2 \psi^*}{\partial x^2} - \frac{1}{i\hbar} V \psi^* \quad \text{+ complex conjugate}$$

$$\Rightarrow \frac{\partial \rho}{\partial t} = -\frac{\hbar}{2mi} \left(\psi^* \frac{\partial^2 \psi}{\partial x^2} - \psi \frac{\partial^2 \psi^*}{\partial x^2} \right) + \frac{1}{i\hbar} \underbrace{(\psi^* V \psi - \psi V \psi^*)}_{\emptyset}$$

$$\frac{\partial \rho}{\partial t} = -\frac{\hbar}{2mi} \left(\psi^* \frac{\partial^2 \psi}{\partial x^2} - \psi \frac{\partial^2 \psi^*}{\partial x^2} \right)$$

$$\frac{\partial j}{\partial x} = \frac{\hbar}{2mi} \left[\frac{\partial \psi^*}{\partial x} \frac{\partial \psi}{\partial x} + \psi^* \frac{\partial^2 \psi}{\partial x^2} - \frac{\partial^2 \psi^*}{\partial x^2} \psi - \frac{\partial \psi^*}{\partial x} \frac{\partial \psi}{\partial x} \right]$$

$$= \frac{\hbar}{2mi} \left(\psi^* \frac{\partial^2 \psi}{\partial x^2} - \psi \frac{\partial^2 \psi^*}{\partial x^2} \right) = -\frac{\partial \rho}{\partial t}$$

$$\Rightarrow \frac{\partial \rho}{\partial t} + \frac{\partial j}{\partial x} = 0 \quad \text{same as charge conservation}$$

or mass continuity for an ideal
incompressible fluid!

$$4. \quad \psi = \sqrt{\frac{m\omega_0}{\pi\hbar}} e^{-\alpha^2 x^2} \quad |\psi|^2 = \sqrt{\frac{m\omega_0}{\pi\hbar}} e^{-2\alpha^2 x^2}$$

$$\langle x \rangle = \int_{-\infty}^{\infty} x |\psi|^2 dx = \int_{-\infty}^{\infty} \sqrt{\frac{m\omega_0}{\pi\hbar}} x e^{-2\alpha^2 x^2} dx = 0 \quad \text{oddeven function}$$

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2 |\psi|^2 dx = \int_{-\infty}^{\infty} \sqrt{\frac{m\omega_0}{\pi\hbar}} x^2 e^{-2\alpha^2 x^2} dx = \frac{1}{2} \sqrt{\frac{m\omega_0}{\pi\hbar}} \underbrace{\int_{-\infty}^{\infty} x^2 e^{-2\alpha^2 x^2} dx}_{\text{lookup}} = \frac{1}{4} \sqrt{\frac{m\omega_0}{2\hbar\alpha^6}}$$

$$\Delta x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \frac{1}{2} \sqrt{\frac{m\omega_0}{2\hbar\alpha^6}}$$

using $\alpha^2 = \frac{m\omega_0}{2\hbar}$ (problem 8)

$$\langle x^2 \rangle = \frac{\hbar}{2m\omega_0}, \quad \Delta x = \sqrt{\frac{\hbar}{2m\omega_0}}$$

$$5. \quad p_{\min} = \Delta p_{\min} = \frac{\hbar}{2\Delta x} = \frac{\hbar}{2x_{\min}}$$

$$E = \frac{\hbar^2}{4x_{\min}^2} \cdot \frac{1}{2m} + \frac{1}{2} k x_{\min}^2 = \frac{\hbar^2}{8m x_{\min}^2} + \frac{1}{2} k x_{\min}^2 = \frac{\hbar^2}{8m x_{\min}^2} + \frac{1}{2} m\omega_0^2 x_{\min}^2$$

$$k = \omega_0^2 m$$

$$\frac{\partial E}{\partial x_{\min}} = \frac{-\hbar^2}{4m x_{\min}^3} + m\omega_0^2 x_{\min} = 0 \Rightarrow 4m^2 \omega_0^2 x_{\min}^4 = \hbar^2$$

$$\Rightarrow x_{\min}^2 = \frac{\hbar}{2m\omega_0}$$

$$\Rightarrow E_{\min} = \frac{\hbar^2}{8m} \frac{2m\omega_0}{\hbar} + \frac{1}{2} m\omega_0^2 \frac{\hbar}{2m\omega_0} = \frac{\hbar\omega_0}{4} + \frac{\hbar\omega_0}{4} = \frac{1}{2} \hbar\omega_0 = \frac{1}{2} \hbar f \checkmark$$

6. particle confined to a region with $L = 10^{-14} \text{ m}$
 in 1D, particle in a box $E = \frac{\hbar^2}{8mL^2} \approx 2 \text{ MeV}$

7. $|\psi|^2 = A^2$

a) $1 = \int_{-\infty}^{\infty} P(x) dx = \int_{-\infty}^{\infty} |\psi|^2 dx = \int_{-b}^{5b} A^2 dx = 6bA^2 \Rightarrow A^2 = \frac{1}{6b}$
 $|\psi|^2 = 0$ outside $[-b, 5b]$

b) $P(\text{in } [0, b]) = \int_0^b |\psi|^2 dx = \int_0^b A^2 dx = A^2 b = \frac{1}{6}$

ψ is constant over interval of width $6b$, so probability of finding it in an interval of width b should be $\frac{1}{6}$

c) $\langle x \rangle = \int_{-b}^{5b} A^2 x dx = \frac{x^2}{2} \Big|_{-b}^{5b} A^2 = \left(\frac{25b^2 - b^2}{2} \right) A^2 = 12b^2 \cdot \frac{1}{6b} = 2b$

$\langle x^2 \rangle = \int_{-b}^{5b} A^2 x^2 dx = \frac{x^3}{3} \Big|_{-b}^{5b} A^2 = \frac{125b^3 - (-b^3)}{3} A^2 = \frac{126}{3} b^3 A^2$

$\langle x^2 \rangle = 42b^3 A^2 = 7b^2$

d) $\Delta x = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{7b^2 - 4b^2} = b\sqrt{3}$

8. $\psi = a e^{-\alpha^2 x^2}$ normalize to find a

$$1 = \int_{-\infty}^{\infty} |\psi|^2 dx = a^2 \int_{-\infty}^{\infty} e^{-2\alpha^2 x^2} dx = \frac{\sqrt{\pi}}{\alpha} a^2 \Rightarrow \alpha^2 = \pi a^4 \text{ need more info...}$$

use Schrödinger equ.

$$\frac{\partial \psi}{\partial x} = -2\alpha^2 x a e^{-\alpha^2 x^2}$$

$$\frac{\partial^2 \psi}{\partial x^2} = -2\alpha^2 a e^{-\alpha^2 x^2} + 4\alpha^4 x^2 a e^{-\alpha^2 x^2}$$

$$= (4\alpha^4 x^2 - 2\alpha^2) \psi$$

$$\Rightarrow \underline{E\psi} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + \frac{1}{2} m \omega_0^2 x^2 \psi = -\frac{\hbar^2}{2m} (4\alpha^4 x^2 - 2\alpha^2) \psi + \frac{1}{2} m \omega_0^2 x^2 \psi$$

can only work for all x if constant terms and x^2 terms multiplying ψ are separately equal

$$\Rightarrow E = -\frac{\hbar^2}{2m} (-2\alpha^2) = \frac{\hbar^2 \alpha^2}{m}$$

$$\Rightarrow -\frac{\hbar^2}{2m} (4\alpha^4) x^2 + \frac{1}{2} m \omega_0^2 x^2 = 0$$

$$\frac{2\hbar^2 \alpha^4}{m} = \frac{1}{2} m \omega_0^2, \quad \alpha^4 = \frac{m^2 \omega_0^2}{4\hbar^2} \quad (\text{for all } x \neq 0)$$

Combine: $E = \frac{\hbar^2 \alpha^2}{m} = \frac{\hbar^2}{m} \frac{m \omega_0}{2\hbar} = \frac{1}{2} \hbar \omega_0$ (same as problem 5)

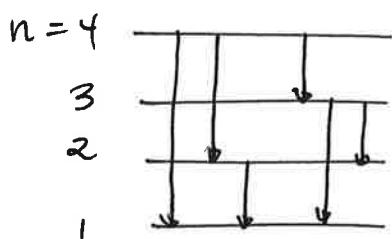
$$\alpha = \sqrt[4]{\frac{m^2 \omega_0^2}{4\hbar^2}} = \sqrt{\frac{m \omega_0}{2\hbar}}$$

$$a^2 = \alpha \sqrt{\pi}, \quad a = \sqrt[4]{\pi} \sqrt{\alpha} = \sqrt[4]{\frac{\pi m \omega_0}{2\hbar}} \quad \checkmark$$

9. $L = 0.17 \text{ nm} = 1.7 \times 10^{-10} \text{ m}$, $n = 4$

$$E_1 = \frac{h^2}{8mL^2} = 13.03 \text{ eV} \quad E_n = n^2 E_1$$

starting from $n=4$, there are $\binom{4}{2} = 6$ possible transitions



transition	energy
$4 \rightarrow 1$	$E_4 - E_1 = (4^2 - 1^2)E_1 = 195.4 \text{ eV}$
$4 \rightarrow 2$	$E_4 - E_2 = (4^2 - 2^2)E_1 = 156.3 \text{ eV}$
$4 \rightarrow 3$	$E_4 - E_3 = (4^2 - 3^2)E_1 = 91.2 \text{ eV}$
$3 \rightarrow 2$	$E_3 - E_2 = (3^2 - 2^2)E_1 = 65.1 \text{ eV}$
$3 \rightarrow 1$	$E_3 - E_1 = (3^2 - 1^2)E_1 = 104.2 \text{ eV}$
$2 \rightarrow 1$	$E_2 - E_1 = (2^2 - 1^2)E_1 = 39.1 \text{ eV}$

photon emitted has energy equal to the difference in energy between initial and final levels. going from level n to m ,

$$\Delta E = hf = E_n - E_m = \frac{n^2 h^2}{8mL^2} - \frac{m^2 h^2}{8mL^2} = (n^2 - m^2) \frac{h^2}{8mL^2} = (n^2 - m^2) E_1$$

10.

$$\Delta x \Delta p \geq \hbar/2 \quad p = mv \quad \Delta p = m \Delta v \quad \text{if mass is known}$$

$$\hbar = h/2\pi = 1.05 \text{ J}\cdot\text{s} \quad \Delta x = 1 \text{ m}$$

$$\Rightarrow m \Delta x \Delta v \geq \hbar/2 \quad \text{or} \quad \Delta v \geq \frac{\hbar}{2m \Delta x} \approx 0.29 \text{ m/s}$$