PH253 Lecture 15: Schrödinger's equation. Still. 1-D potentials

P. LeClair

Department of Physics & Astronomy The University of Alabama

Spring 2020



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Image of the day

High-resolution electron microscope image of a $Co_{1.5}Ti_{0.5}FeGe$ alloy sample. The bright dots are individual atoms. The inset is a Fourier transform of the image, indicating the hexagonal symmetry.

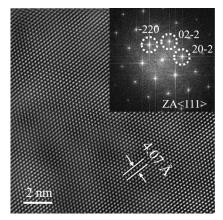


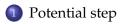
Figure: UA Physics / LeClair group / Phys. Rev. Materials 3, 114406 (2019)

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Outline



2 Particle in a box





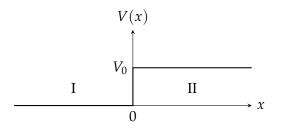
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Potential step from last time





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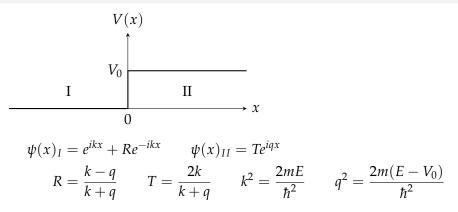
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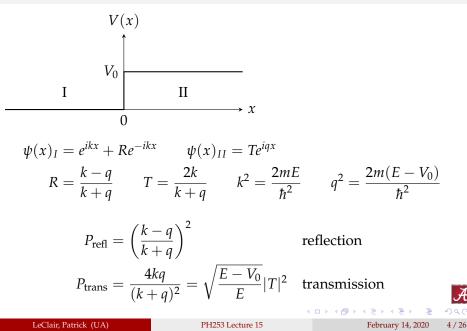
Potential step from last time





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Potential step from last time



$$P_{\text{refl}} = \left(\frac{k-q}{k+q}\right)^2 \qquad \text{reflection}$$
$$P_{\text{trans}} = \frac{4kq}{(k+q)^2} = \sqrt{\frac{E-V_0}{E}}|T|^2 \quad \text{transmission}$$



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① Degree of transmission depends on k, q, i.e., E compared to V_o



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High energy: *k* ~ *q* gives P_{refl} ~ 0, P_{trans} ~ 1



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- Solution Low energy: $k \gg q$ gives significant P_{refl}



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- **1** Degree of transmission depends on k, q, i.e., E compared to V_o
- **2** High energy: $k \sim q$ gives $P_{\text{refl}} \sim 0$, $P_{\text{trans}} \sim 1$
- **(a)** Low energy: $k \gg q$ gives significant P_{refl}
- Check: k = q, $P_{\text{trans}} = 1$ there is no step!

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- So Check: q = 0, $P_{\text{trans}} = 0$ zero energy in II!

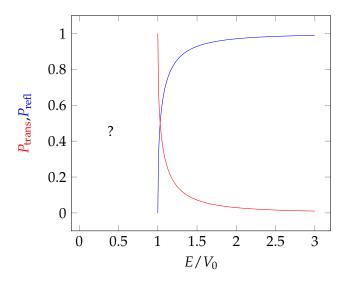
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- Check: k = q, $P_{\text{trans}} = 1$ there is no step!
- So Check: q = 0, $P_{\text{trans}} = 0$ zero energy in II!
- $E = V_0$? $P_{\text{trans}} = 0$, perfect reflection



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• Classical: go over step, slow down to conserve *E*, never reflect



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- Quantum: chance of reflection, even if *E* high enough.



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- Solution If $E < V_0$, *q* is imaginary oscillation becomes decay!



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$$\psi_{II} = Te^{-|q|x}$$
 $T = \frac{2k}{k+i|q|} \neq 0$ $(E < V_0)$



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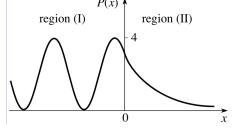


Figure: http://www.met.reading.ac.uk/pplato2/h-flap/phys11_1.html



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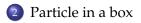
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O Potential zero inside, infinite outside



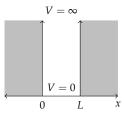
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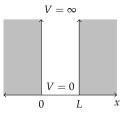
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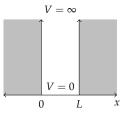
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- Can also imagine charged plates + e^-

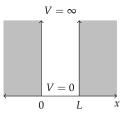


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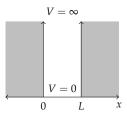


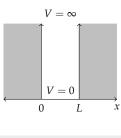
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- Solution $\psi(0) = \psi(L) = 0!$





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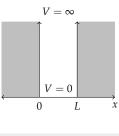
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1 Wave can't penetrate boundary, so ψ goes to zero there





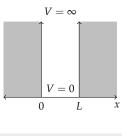
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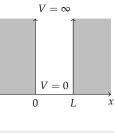
- **1** Wave can't penetrate boundary, so ψ goes to zero there
- Spoiler alert: standing waves





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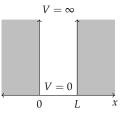
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$$E = \frac{p^2}{2m} = \frac{\hbar^2 k^2}{2m}$$





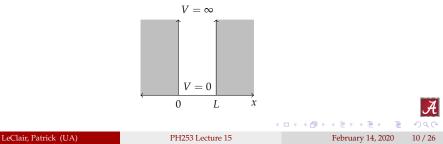
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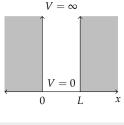
Now set up Schrödinger's equation in box + boundary conditions.



Boundary conditions are key

Schrödinger equation with V = 0:

$$\frac{d^2\psi}{dx^2} = -\left(\frac{2mE}{\hbar^2}\right)\psi = -k^2\psi \qquad \left(\text{with } k^2 = \frac{2mE}{\hbar^2}\right)$$



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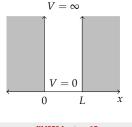
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We already know the solution





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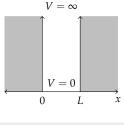
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We already know the solution

$$\psi(x) = C_1 e^{ikx} + C_2 e^{-ikx}$$





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$$\psi(x) = C_1 e^{ikx} + C_2 e^{-ikx}$$

How to find constants? Boundary conditions! $\psi(0) = \psi(L) = 0$ $V = \infty$

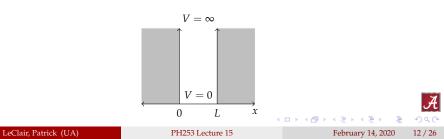
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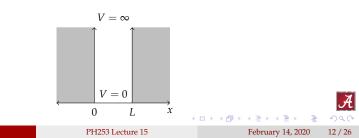


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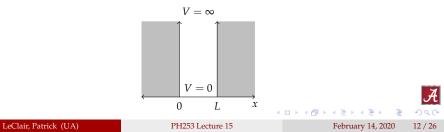


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Note
$$e^{ikx} - e^{-ikx} \propto \sin kx \implies \psi(x) = C \sin kx$$
.

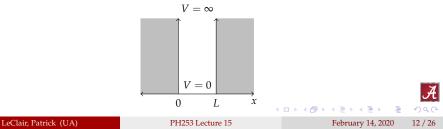


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$$\psi(x) = C_1 e^{ikx} - C_1 e^{-ikx}$$

Note $e^{ikx} - e^{-ikx} \propto \sin kx \implies \psi(x) = C \sin kx$. Recall: bound states have ψ purely real



$$\psi(L) = C\sin kL = 0$$



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$$\psi(L) = C\sin kL = 0$$

• Only true when $kL = n\pi$, $n = \{1, 2, 3, ...\}$.



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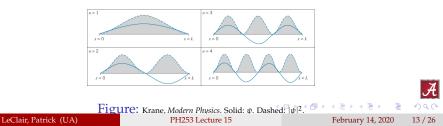
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- *n* is like a harmonic, how many half λ 's fit in box



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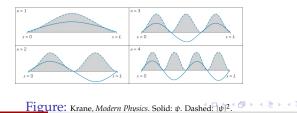
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- *n* is like a harmonic, how many half λ 's fit in box
- $k_n = \frac{n\pi}{L}$, or $n\frac{\lambda}{2} = L$. Implies discrete energies too...



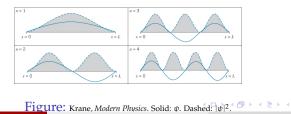
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- Basically: modes of a string. *C* might depend on *n*?



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$$E_n = \frac{\hbar^2 k_n^2}{2m} = \frac{\hbar^2 \pi^2 n^2}{2mL^2} = \frac{n^2 h^2}{8mL^2}$$

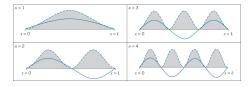
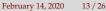


Figure: Krane, Modern Physics. Solid: ψ . Dashed: $|\psi|^2$. PH253 Lecture 15



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$$\psi(x) = \begin{cases} C_n \sin\left(\frac{n\pi x}{L}\right) & 0 \le x \ge L\\ 0 & x < 0, x > L \end{cases}$$

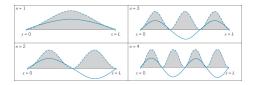


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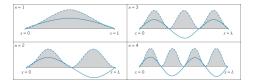


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Big picture: *confinement* leads to quantized energy levels.

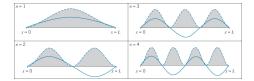


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- **(**) Since *V* is independent of time, time variation is $e^{-iE_nt/\hbar}$
- **2** Big picture: *confinement* leads to quantized energy levels.
- Where we went wrong with free particle: no boundary conditions

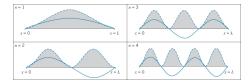


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- **2** Big picture: *confinement* leads to quantized energy levels.
- 9 Where we went wrong with free particle: no boundary conditions
- **(9)** What are the C_n ? We have not normalized the wave function ...

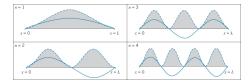


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What are the C_n? We have not normalized the wave function ...
 ≥ Enforce ∫[∞]_{-∞} |ψ(x)|² dx = 1

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$$\psi(x) = \begin{cases} C_n \sin\left(\frac{n\pi x}{L}\right) & 0 \le x \ge L\\ 0 & x < 0, x > L \end{cases}$$

What are the *C_n*? We have not normalized the wave function ...

- 2 Enforce $\int_{-\infty}^{\infty} |\psi(x)|^2 dx = 1$
- Integrand only nonzero from 0 to L though!



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$$\psi(x) = \begin{cases} C_n \sin\left(\frac{n\pi x}{L}\right) & 0 \le x \ge L\\ 0 & x < 0, x > L \end{cases}$$

• What are the *C_n*? We have not normalized the wave function ...

2 Enforce $\int_{-\infty}^{\infty} |\psi(x)|^2 dx = 1$

Integrand only nonzero from 0 to L though!

$$1 = \int_{0}^{L} |\psi(x)|^{2} dx = \int_{0}^{L} C_{n}^{2} \sin^{2}\left(\frac{n\pi x}{L}\right) dx$$
$$= C_{n}^{2} \left[\frac{x}{2} - \frac{4L}{n\pi} \sin\left(2\frac{n\pi x}{L}\right)\right] \Big|_{0}^{L} = C_{n}^{2} \frac{L}{2}$$

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•
$$C_n^2(L/2) = 1$$
, so $C_n = \sqrt{2/L}$ for all *n*



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- So far we can get energy from the Schrödinger equation
- ② We figured out the momentum operator last time, $p = -i\hbar \frac{\partial}{\partial x}$
- What about position, and how to use these new operators?

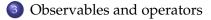


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Outline









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• We can calculate *expectation values* for any observable



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- We can calculate *expectation values* for any observable
- With many measurements, what would you expect on average?



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- **③** Generic form for observable *A* is $\langle A \rangle = \int \psi^* A_{op} \psi \, dx$
- So The position operator is simple, it is just *x*

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Observables

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$$\langle x \rangle = \int_{-\infty}^{+\infty} \psi^* x \psi \, dx = \int_{-\infty}^{+\infty} x |\psi(x)|^2 \, dx = \int_{-\infty}^{+\infty} x P(x), dx$$

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Mostly stats – uncertainty from probability distribution.



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Mostly stats – uncertainty from probability distribution.
 (Δx)² = (x²) - (x)²



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- Mostly stats uncertainty from probability distribution.
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- **③** The same Δx from the uncertainty principle



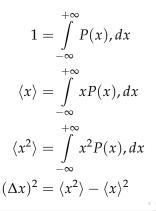
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• Potential energy is easy: operator is just V(x).



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Potential energy is easy: operator is just V(x). ⟨U⟩ = ∫ V(x)|ψ(x)|² dx



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• Potential energy is easy: operator is just V(x).

$$\ 0 \ \ \langle U \rangle = \int V(x) |\psi(x)|^2 \, dx$$

Momentum is less nice. For operators, order matters in general

$$\mathbf{p}_{\text{oper}} = \frac{\hbar}{i} \frac{\partial}{\partial x} \qquad \mathbf{p}_{\text{oper}}^2 = -\hbar^2 \frac{\partial^2}{\partial x^2}$$
$$\langle p \rangle = \int_{-\infty}^{+\infty} \psi^*(x) \frac{\hbar}{i} \frac{\partial}{\partial x} \psi(x) \, dx$$
$$\langle p^2 \rangle = \int_{-\infty}^{+\infty} \psi^*(x) \left(-\hbar^2 \frac{\partial^2}{\partial x^2} \right) \psi(x) \, dx$$
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But! Can now formalize uncertainty principle & actually calculate.





• What is $\langle x \rangle$ for the particle in a box?



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- What is $\langle x \rangle$ for the particle in a box?
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$$\begin{aligned} \langle x \rangle &= \int_{-\infty}^{+\infty} x |\psi(x)|^2 \, dx = \int_{0}^{L} x \frac{2}{L} \sin^2 \frac{n\pi x}{L} \, dx \qquad \text{let } u = n\pi x/L \\ &= \frac{2}{L} \frac{L}{n\pi} \int_{0}^{n\pi} \frac{L}{n\pi} u \sin^2 u \, du = \frac{2L}{n^2 \pi^2} \left[\frac{u^2}{4} - \frac{1}{4} u \sin 2u - \frac{1}{8} \cos 2u \right]_{0}^{n\pi} \\ &= \frac{2L}{n^2 \pi^2} \left[\frac{n^2 \pi^2}{4} \right] = \frac{L}{2} \end{aligned}$$

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 $\sin 2u$, $\cos 2u$ terms are zero or cancel. Result as expected.

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• Now we need $\langle x^2 \rangle$. Remember $u = n\pi x/L$



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sin 2*u*, cos 2*u* terms are zero or cancel.
 Then Δx = √L²/3 − (L/2)² = L/2√3 ≈ L/3.46

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 \bigcirc sin 2*u*, cos 2*u* terms are zero or cancel.

2 Then $\Delta x = \sqrt{L^2/3 - (L/2)^2} = L/2\sqrt{3} \approx L/3.46$

Solution Measurement? $x_{\text{best}} = \langle x \rangle \pm \langle x^2 \rangle = (0.500 \pm 0.289)L - \text{broad}$



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I How about momentum?



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- On't need math. Just as much time backwards as forward!



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- **(**) $\langle p \rangle = 0$ by symmetry, just as $\langle x \rangle = L/2$ must be right



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Average momentum is zero, uncertainty/spread?
Need (p²) for that

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- Average momentum is zero, uncertainty/spread?
- **2** Need $\langle p^2 \rangle$ for that
- So We know $p^2 = 2mE$ for a free particle, must be true in box too ...

() Since
$$p^2 = 2mE$$
, then $\langle p^2 \rangle = 2mE_n$ – know *E* already



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In any case:

$$\Delta p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \sqrt{2mE_n - 0} = \frac{n\pi\hbar}{L}$$



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1 Put it together: we know Δx and Δp now.



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- Put it together: we know Δx and Δp now.
- Ooes the uncertainty principle hold?



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$$\Delta x \Delta p = \left(\frac{L}{2\sqrt{3}}\right) \left(\frac{n\pi\hbar}{L}\right) = \frac{n\pi\hbar}{2\sqrt{3}} > \frac{\hbar}{2}$$



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• Numerical, $\Delta x \Delta p \approx 0.9n\hbar > 0.5\hbar$

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- **2** Well above 0.5*h* limit, by factor $n\pi/\sqrt{3} \approx 1.81n$

- Put it together: we know Δx and Δp now.
- 2 Does the uncertainty principle hold?

$$\Delta x \Delta p = \left(\frac{L}{2\sqrt{3}}\right) \left(\frac{n\pi\hbar}{L}\right) = \frac{n\pi\hbar}{2\sqrt{3}} > \frac{\hbar}{2}$$

- Numerical, $\Delta x \Delta p \approx 0.9n\hbar > 0.5\hbar$
- **2** Well above 0.5*h* limit, by factor $n\pi/\sqrt{3} \approx 1.81n$
- Solution Not unreasonable: uncertainty up as n (and E_n) increase

Find potential, by region if necessary



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- Find potential, by region if necessary
- Write down and solve Schrödinger's equation for each region



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- § Find overall constants by normalization

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- Find potential, by region if necessary
- Write down and solve Schrödinger's equation for each region
- Senforce any boundary conditions you know
- **9** Enforce continuity of ψ and its derivatives at boundaries
- § Find overall constants by normalization
- O Up next: modeling atoms

