

PH253 Lecture 4: relativity

Lorentz Transformations, Dynamics

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Outline

- 1 Lorentz Transformations
- 2 Doppler shift
- 3 Momentum
- 4 Energy



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Lorentz transformations between coordinate systems

$$x' = \gamma(x - vt) \quad \text{or} \quad x = \gamma(x' + vt') \quad (1)$$

$$y' = y \quad (2)$$

$$z' = z \quad (3)$$

$$t' = \gamma\left(t - \frac{vx}{c^2}\right) \quad \text{or} \quad t = \gamma\left(t' + \frac{vx'}{c^2}\right) \quad (4)$$

Here (x, y, z, t) is the position and time of an event as measured by an observer in O stationary to it. A second observer in O' , moving at velocity v along the x axis, measures the same event to be at position and time (x', y', z', t') .



Velocity Addition

We have an observer in a frame O , and a second observer in another frame O' who are moving relative to each other at a velocity v . Both observers measure the velocity of another object in their own frames (v_{obj} and v'_{obj}). We can relate the velocities measured in the different frames as follows:

$$v_{\text{obj}} = \frac{v + v'_{\text{obj}}}{1 + \frac{vv'_{\text{obj}}}{c^2}} \quad v'_{\text{obj}} = \frac{v_{\text{obj}} - v}{1 - \frac{vv_{\text{obj}}}{c^2}} \quad (5)$$

Again, v_{obj} is the object's velocity as measured from the O reference frame, and v'_{obj} is its velocity as measured from the O' reference frame.



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Decide if you should add or subtract velocities, then pick the formula.



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$$\frac{dy'}{dt} = \frac{dy}{dt} = v_y \quad \text{no length contraction} \quad (7)$$

$$\frac{dt'}{dt} = \frac{d}{dt} \left(\gamma \left(t - \frac{vx}{c^2} \right) \right) = \gamma \left(1 - \frac{vv_x}{c^2} \right) \quad (8)$$



Velocity Transformations?

Put it together:

$$v'_y = \frac{v_y}{\gamma \left(1 - \frac{vv_x}{c^2}\right)} \quad (9)$$



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Transverse velocities do change! Similar expression along z .

This is a homework problem.



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Expect this for light, but $v_{\text{light}} \gg v_{\text{sound}}$... harder to observe



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- *And*, crests of waves are coming faster: $v'_w = v_w + v_o$



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Moving *toward* source at v_o shortens wavelength λ toward blue

Similarly, moving *away* source at v_o lengthens wavelength λ toward red ($v_o < 0$)



Doppler shift?

Red shift: ratio of shifted wavelength to what a stationary observer sees:

$$z = \frac{\lambda'}{\lambda} = \sqrt{\frac{c + v_o}{c - v_o}} \quad (12)$$

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Relativity principle: makes no difference if source or observer is moving.



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$$\text{proper velocity} = \eta = \left(\frac{\text{rate of change of posn. in } \textit{observer's} \text{ frame}}{\text{proper time measured by moving observer}} \right)$$



Momentum

$$\eta = \frac{dl}{dt_p} = \frac{dl/dt}{dt_p/dt} = \frac{v}{dt_p/dt} \quad (13)$$

$$\frac{dt_p}{dt} = \frac{d}{dt} (t/\gamma) = \frac{1}{\gamma} \quad (14)$$

$$\Rightarrow \eta = \frac{v}{\sqrt{1 - v^2/c^2}} = \gamma v \quad (15)$$



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What about mass? Modern view: *invariant* - all agree on rest mass



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Energy

The work-energy theorem relates the change in kinetic energy of a particle to the work done on it by an external force: $\Delta K = W = \int F dx$. Writing Newton's second law as $F = dp/dt$, show that $W = \int v dp$ and integrate by parts to obtain the result

$$K = \frac{mc^2}{\sqrt{1 - v^2/c^2}} - mc^2$$



Energy

We are concerned with finding the kinetic energy of a particle accelerated from rest, starting with zero kinetic energy, so we may simply write $K = W = \int F dx$ and drop the Δ . Using what we have, and noting $v = dx/dt$:

$$K = \int_{v=0}^{v=v} F dx = \int_0^v \frac{dp}{dt} dx = \int_0^v dp \frac{dx}{dt} = \int_0^v v dp \quad (16)$$



Energy

We can integrate by parts. Let $f = dp$ and $g' = dx/dt$, and recall integration by parts gives $\int fg' = fg - \int f'g$.

$$K = \int_0^v v dp = pv \Big|_0^v - \int_0^v p dv = pv - \int_0^v p dv \quad (17)$$



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Noting $p = mv / \sqrt{1 - v^2/c^2}$,

$$K = \frac{mv^2}{\sqrt{1 - v^2/c^2}} - \int_0^v \frac{mv}{\sqrt{1 - v^2/c^2}} dv \quad (18)$$

$$= \frac{mv^2}{\sqrt{1 - v^2/c^2}} + mc^2 \sqrt{1 - v^2/c^2} \Big|_0^v \quad (19)$$



Energy

Continuing:

$$K = \frac{mv^2}{\sqrt{1 - v^2/c^2}} + mc^2 \sqrt{1 - v^2/c^2} - mc^2 \quad (20)$$

$$= \frac{mv^2 + mc^2 (1 - v^2/c^2)}{\sqrt{1 - v^2/c^2}} - mc^2 \quad (21)$$

$$= \frac{mc^2}{\sqrt{1 - v^2/c^2}} - mc^2 = \left(\frac{1}{\sqrt{1 - v^2/c^2}} - 1 \right) mc^2 \quad (22)$$

$$\implies K = (\gamma - 1) mc^2 \quad (23)$$



Energy

Low v ? For $v \ll c$, $\gamma = (1 - v^2/c^2)^{-1/2} \approx 1 + \frac{v^2}{c^2}$



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Plug this in: $K = (\gamma - 1) mc^2 \approx \left(1 + \frac{v^2}{c^2} - 1\right) mc^2 = \frac{1}{2}mv^2$



- Now if $K = (\gamma - 1) mc^2$, as $v \rightarrow c$ then $K \rightarrow \infty$.



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- Takes infinite energy to reach light speed, not happening
- Only massless objects (light) can travel at c
- What about second term, independent of speed?!?



- $K = (\gamma - 1) mc^2 = \gamma mc^2 - mc^2$



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- $K = (\gamma - 1) mc^2 = \gamma mc^2 - mc^2$
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- Or, $E_{\text{tot}} = K + mc^2 = \gamma mc^2$ and $E_{\text{rest}} = mc^2$
- Bodies have *rest energy* associated with mass – mass-energy equivalent!



Relativistic energy-momentum equations

Can go further ...

$$E_{\text{tot}}^2 - (pc)^2 = (mc^2)^2 \quad (24)$$

$$E_{\text{tot}}v = pc^2 \quad (25)$$

Better form: relativistic E – p relationship



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$$\text{KE} = \sqrt{p^2c^2 + m^2c^4} - mc^2 \quad \text{and} \quad p = \sqrt{\frac{E^2}{c^2} - m^2c^4} \quad (26)$$



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c is really a speed limit for any massive particle

