## Lattice distortions in cubic crystals

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## TETRAGONAL DISTORTION

For a cubic crystal, we know that for instance the $\{100\}$ family of planes all give the same reflections - (100), (010), and (001) all give a diffraction peak at the same angle. Now consider a small distortion along the $c$ axis. That makes (001) not equivalent to (100) and (010), and the single diffraction peak now becomes a doublet: a peak of relative intensity 2 for the (100) and (010) overlapping reflections, and a peak of relative intensity 1 for the ( 001 ) reflection. Below we derive an expression for the ratio of the lattice constants $a / c$ presuming that it is close to unity, i.e., small distortions.

For a tetragonal system, we have

$$
\begin{align*}
\sin ^{2} \theta & =\frac{\lambda^{2}}{4 a^{2}}\left(h^{2}+k^{2}\right)+\frac{\lambda^{2}}{4 c^{2}} l^{2}  \tag{1}\\
\Longrightarrow \quad \frac{4 \sin ^{2} \theta}{\lambda^{2}} & =\frac{h^{2}+k^{2}}{a^{2}}+\frac{l^{2}}{c^{2}} \tag{2}
\end{align*}
$$

For the sake of concreteness, consider a reflection (hhl) occurring at angle $\theta_{1}$ and reflections (lhh) and (hlh) occuring at angle $\theta_{2}$. The difference in $\sin ^{2}$ values is then

$$
\begin{equation*}
\sin ^{2} \theta_{1}-\sin ^{2} \theta_{2}=\frac{\lambda^{2}}{4}\left(\frac{h^{2}+h^{2}}{a^{2}}+\frac{l^{2}}{c^{2}}-\frac{l^{2}+h^{2}}{a^{2}}-\frac{h^{2}}{c^{2}}\right)=\frac{\lambda^{2}\left(h^{2}-l^{2}\right)}{4}\left(\frac{1}{a^{2}}-\frac{1}{c^{2}}\right)=\frac{\lambda^{2}\left(h^{2}-l^{2}\right)}{4 a^{2}}\left(1-\frac{a^{2}}{c^{2}}\right) \tag{3}
\end{equation*}
$$

We can simplify the left-hand side. Let $\theta_{o}=\left(\theta_{1}+\theta_{2}\right) / 2$ and $\delta=\theta_{1}-\theta_{2}$, so $\theta_{1}=\theta_{o}+\delta / 2$ and $\theta_{2}=\theta_{o}-\delta / 2$. This make $\theta_{o}$ the original peak position before distortion, and $\delta$ is the peak splitting after the distortion. Then:

$$
\begin{align*}
\sin ^{2} \theta_{1}-\sin ^{2} \theta_{2}= & \sin ^{2}\left(\theta_{o}+\frac{\delta}{2}\right)-\sin ^{2}\left(\theta_{o}-\frac{\delta}{2}\right)  \tag{4}\\
= & \left(\sin \theta_{o} \cos \frac{\delta}{2}+\cos \theta_{o} \sin \frac{\delta}{2}\right)^{2}-\left(\sin \theta_{o} \cos \frac{\delta}{2}-\cos \theta_{o} \sin \frac{\delta}{2}\right)^{2}  \tag{5}\\
= & \sin ^{2} \theta_{o} \cos ^{2} \frac{\delta}{2}+2 \sin \theta_{o} \cos \theta_{o} \sin \frac{\delta}{2} \cos \frac{\delta}{2}+\cos ^{2} \theta_{o} \sin ^{2} \frac{\delta}{2}  \tag{6}\\
& -\sin ^{2} \theta_{o} \cos ^{2} \frac{\delta}{2}+2 \sin \theta_{o} \cos \theta_{o} \sin \frac{\delta}{2} \cos \frac{\delta}{2}-\cos ^{2} \theta_{o} \sin ^{2} \frac{\delta}{2}  \tag{7}\\
= & 4 \sin \theta_{o} \cos \theta_{o} \sin \frac{\delta}{2} \cos \frac{\delta}{2}=2 \sin 2 \theta_{o} \sin \frac{\delta}{2} \cos \frac{\delta}{2} \tag{8}
\end{align*}
$$

Since $\delta$ is small, $\cos \frac{\delta}{2} \approx 1$ and $\sin \frac{\delta}{2} \approx \frac{\delta}{2}$, and

$$
\begin{equation*}
\sin ^{2} \theta_{1}-\sin ^{2} \theta_{2} \approx \delta \sin 2 \theta_{o} \tag{9}
\end{equation*}
$$

Using Eq. 2,

$$
\begin{align*}
\delta \sin 2 \theta_{o} & \approx \frac{\lambda^{2}\left(h^{2}-l^{2}\right)}{4 a^{2}}\left(1-\frac{a^{2}}{c^{2}}\right)  \tag{10}\\
\frac{a}{c} & \approx \sqrt{1-\frac{4 a^{2} \delta \sin 2 \theta_{o}}{\lambda^{2}\left(h^{2}-l^{2}\right)}} \tag{11}
\end{align*}
$$

Thus, once the peaks are indexed, from the peak splitting $\delta$ and the average angle $2 \theta_{o}$ we can find the $a / c$ ratio.

## RHOMBOHEDRAL DISTORTION

For a rhombohedral crystal with an angle $\alpha$ between axes, we have

$$
\begin{equation*}
\frac{1}{d^{2}}=\frac{\left(h^{2}+k^{2}+l^{2}\right) \sin ^{2} \alpha+2(h k+k l+h l)\left(\cos ^{2} \alpha-\cos \alpha\right)}{a^{2}\left(1-3 \cos ^{2} \alpha+2 \cos ^{3} \alpha\right)} \tag{12}
\end{equation*}
$$

Here we assume we have a nearly cubic crystal with a small rhombohedral distortion, such that the lattice angle is $\alpha=\frac{\pi}{2}+\delta$, where $\delta$ is the distortion. This corresponds to a stretch (or compression) along the body-diagonal of the cubic unit cell. This leaves the 100 cubic peak unsplit, but splits the 111 cubic reflection into two peaks (111 and $11 \overline{1})$ since there are two different body-diagonal distances in the rhombohedral unit cell.

In this case, for small $\delta$,

$$
\begin{align*}
& \sin \alpha=\sin \left(\frac{\pi}{2}+\delta\right)=\cos \delta \approx 1-\frac{\delta^{2}}{2}  \tag{13}\\
& \cos \alpha=\cos \left(\frac{\pi}{2}+\delta\right)=-\sin \delta \approx-\delta \tag{14}
\end{align*}
$$

This gives

$$
\begin{equation*}
\frac{1}{d^{2}}=\frac{\left(h^{2}+k^{2}+l^{2}\right)\left(1-\frac{\delta^{2}}{2}\right)^{2}+2(h k+k l+h l)\left(\delta^{2}+\delta\right)}{a^{2}\left(1-3 \delta^{2}-2 \delta^{3}\right)} \tag{15}
\end{equation*}
$$

Expanding, and neglecting terms of order $\delta^{4}$ and higher (so $\left(1-\delta^{2} / 2\right)^{2} \approx 1-\delta^{2}$ ),

$$
\begin{align*}
\frac{1}{d^{2}} & =\frac{\left(h^{2}+k^{2}+l^{2}\right)\left(1-\delta^{2}\right)+2 \delta(h k+k l+h l)(1+\delta)}{a^{2}\left(1-3 \delta^{2}-2 \delta^{3}\right)}  \tag{16}\\
& =\frac{\left(h^{2}+k^{2}+l^{2}\right)(1-\delta)(1+\delta)+2 \delta(h k+k l+h l)(1+\delta)}{a^{2}(1+\delta)\left(1-\delta-2 \delta^{2}\right)}  \tag{17}\\
& =\frac{\left(h^{2}+k^{2}+l^{2}\right)(1-\delta)+2 \delta(h k+k l+h l)}{a^{2}\left(1-\delta-2 \delta^{2}\right)} \tag{18}
\end{align*}
$$

In terms of the measured reflection angle $\theta$,

$$
\begin{equation*}
\sin ^{2} \theta=\frac{\lambda^{2}}{4 d^{2}}=\frac{\lambda^{2}}{4 a^{2}} \frac{\left(h^{2}+k^{2}+l^{2}\right)(1-\delta)+2 \delta(h k+k l+h l)}{\left(1-\delta-2 \delta^{2}\right)} \tag{20}
\end{equation*}
$$

In order to obtain $\delta$, we make use of the splitting of fundamental lines. For a doublet with peaks at angles $\theta_{1}$ and $\theta_{2}$, the first term in the numerator will be the same for both peaks, and the difference in $\sin ^{2} \theta$ values simplifies. Assuming the lines are $(h k l)$ and ( $\bar{h} k l$ ),

$$
\begin{equation*}
\frac{4 a^{2}}{\lambda^{2}}\left(\sin ^{2} \theta_{1}-\sin ^{2} \theta_{2}\right)=\frac{2 \delta}{1-\delta-2 \delta^{2}}(h k+k l+h l-\bar{h} k-k l-\bar{h} l)=\frac{2 \delta}{1-\delta-2 \delta^{2}}(2 h k+2 h l)=\frac{4 \delta}{1-\delta-2 \delta^{2}}(h k+h l) \tag{21}
\end{equation*}
$$

Since the lattice distortion is expected to be small (verified by the small degree of doublet splitting), $\delta \ll 1$ (in radians) we may simplify the previous expression to:

$$
\begin{equation*}
\frac{\delta}{1-\delta-2 \delta^{2}} \approx \delta \approx \frac{a^{2}\left(\sin ^{2} \theta_{1}-\sin ^{2} \theta_{2}\right)}{\lambda^{2}(h k+h l)} \tag{22}
\end{equation*}
$$

Using the lattice constant $a, \lambda$, and miller indices for each of the double peaks, one may use the difference between the squared sines of the doublet angles to find the degree of rhombohedral distortion $\delta$. This procedure gives little error for the small distortion angles considered, and has the advantage of simplifying analysis a great deal.

