## University of Alabama

Department of Physics and Astronomy
PH 125 / LeClair
Fall 2017

## Exam III Solution

1. A child throws a ball with an initial speed of $8.00 \mathrm{~m} / \mathrm{s}$ at an angle of $40.0^{\circ}$ above the horizontal. The ball leaves her hand 1.00 m above the ground and experience negligible air resistance. (a) How far from where the child is standing does the ball hit the ground? (b) How long is the ball in flight before it hits the ground?

Solution: Let the ball's starting height be $y_{o}$, the initial velocity be $v_{i}$, the launch angle be $\theta$, and the horizontal distance traveled be $d$. If we put the origin at ground level where the child stands, such that the ball starts at coordinates $\left(0, y_{o}\right)$, we can use the $y(x)$ equation to find the distance traveled by finding the $x$ coordinate at which $y=0$.

$$
\begin{align*}
y(x) & =0=y_{o}+x \tan \theta-\frac{g x^{2}}{2 v_{i}^{2} \cos ^{2} \theta}  \tag{1}\\
0 & =-g x^{2}+2 v_{i}^{2} \cos ^{2} \theta \tan \theta x+2 v_{i}^{2} y_{o} \cos ^{2} \theta  \tag{2}\\
0 & =g x^{2}-2 v_{i}^{2} \sin \theta \cos \theta x-2 v_{i}^{2} y_{o} \cos ^{2} \theta \tag{3}
\end{align*}
$$

Solving the quadratic and simplifying,

$$
\begin{align*}
x & =\frac{1}{2 g}\left(2 v_{i}^{2} \sin \theta \cos \theta \pm \sqrt{4 v_{i}^{2} \sin ^{2} \theta \cos ^{2} \theta+8 v_{i}^{2} y_{o} g \cos ^{2} \theta}\right)  \tag{4}\\
& =\frac{v_{i}^{2}}{g}\left(\sin \theta \cos \theta \pm \cos \theta \sqrt{\sin ^{2} \theta+\frac{2 y_{o} g}{v_{i}^{2}}}\right)  \tag{5}\\
& =\frac{v_{i}^{2} \cos \theta}{g}\left(\sin \theta \pm \sqrt{\sin ^{2} \theta+\frac{2 y_{o} g}{v_{i}^{2}}}\right) \approx 7.46,-1.03 \mathrm{~m} \tag{6}
\end{align*}
$$

Clearly the positive solution is the one we seek, so $d=7.46 \mathrm{~m}$. The time spent in the air is then easily found from the $x$ component of the launch velocity:

$$
\begin{equation*}
d=v_{x} t=v_{i} \cos \theta t \quad \Longrightarrow \quad t=\frac{d}{v_{i} \cos \theta} \approx 1.22 \mathrm{~s} \tag{7}
\end{equation*}
$$

2. A uniform solid sphere of mass $M$ and radius $R$ rotates with an angular speed $\omega$ about an axis through its center. A uniform solid cylinder of mass $M$, radius $R$, and length $2 R$ rotates through an axis running through the central axis of the cylinder. What must be the angular speed of the cylinder so it will have the same rotational kinetic energy as the sphere?

Solution: First: we don't need the length of the cylinder at all. All we need to do is equate the rotational kinetic energy of the two, with the sphere rotating with $\omega_{s}$ and the cylinder at $\omega_{c}$.

$$
\begin{align*}
K_{r, \text { cylinder }} & =K_{r, \text { sphere }}  \tag{8}\\
\frac{1}{2} I_{c} \omega_{c}^{2} & =\frac{1}{2} I_{s} \omega_{s}^{2}  \tag{9}\\
\omega_{c} & =\sqrt{\frac{I_{s}}{I_{c}}} \omega_{s}=\sqrt{\frac{2}{\frac{1}{2}}} \omega_{s}=\frac{2}{\sqrt{5}} \omega_{s} \tag{10}
\end{align*}
$$

3. Two blocks are connected by a string that goes over an ideal pulley as shown in the figure. Block A has a mass of 3.00 kg and can slide over a rough plane inclined $30.0^{\circ}$ to the horizontal. The coefficient of kinetic friction between block A and the plane is 0.400 . Block B has a mass of 2.77 kg . What is the acceleration of the blocks?


Solution: We need free-body diagrams for each mass. Let the $x$ axis run up the ramp for mass $A$. We will assume that mass $B$ falls and pulls mass $A$ up the ramp, meaning the acceleration is in the $x$ direction for mass $A$ and downward for mass $B$. For mass $A$ do we need to consider a friction force. Since the rope is presumably taut the entire time of interest, the acceleration is the same for both blocks. For the same reason, the tension applied to both blocks is the same. Newton's second law and geometry will suffice to find the acceleration if we neglect the pulley's rotational inertia. Along the $y$ direction for either mass, the forces must sum to zero, while along the $x$ direction, the forces must give the acceleration for each mass.

$$
\begin{aligned}
& \sum F_{y}=0 \\
& \sum F_{x}=m a_{x}
\end{aligned}
$$

First consider mass $A$. The free body diagram above yields the following, noting that the acceleration will be purely along the $-x$ direction:

$$
\begin{aligned}
& \sum F_{y}=n-m_{a} g \cos \theta=0 \\
& \sum F_{x}=T-f-m_{a} g \sin \theta=m_{a} a_{x}
\end{aligned}
$$

From the first equation, we see $n=m_{a} g \cos \theta$, so $f=\mu_{k} m_{a} g \cos \theta$. For mass $B$,

$$
\sum F_{y}=T-m_{b} g=-m_{b} a \quad \Longrightarrow \quad T=m_{b}(g-a)
$$

Combining,

$$
\begin{align*}
T-f-m_{a} g \sin \theta & =m_{b} g-m_{b} a-\mu_{k} m_{a} g \cos \theta-m_{a} g \sin \theta=m_{a} a  \tag{11}\\
\left(m_{b}+m_{a}\right) a & =g\left(m_{b}-\mu_{k} m_{a} \cos \theta-m_{a} \sin \theta\right)  \tag{12}\\
a & =\left(\frac{m_{b}-\mu_{k} m_{a} \cos \theta-m_{a} \sin \theta}{m_{a}+m_{b}}\right) \approx 0.39 \mathrm{~m} / \mathrm{s}^{2} \tag{13}
\end{align*}
$$

4. A record is dropped vertically onto a freely rotating (undriven) turntable. Frictional forces act to bring the record and turntable to a common angular speed. If the rotational inertia of the record is 0.54 times that of the turntable, what percentage of the initial kinetic energy is lost?

Solution: The record starts at rest with the turntable rotating with velocity $\omega_{i}$. Afterwards, both rotate together with velocity $\omega_{f}$. Dropping the record on the turntable is essentially an inelastic rotational collision, so conservation of angular momentum will relate the two velocities.

$$
\begin{align*}
L_{i} & =L_{f}  \tag{14}\\
I_{t} \omega_{i} & =\left(I_{d}+I_{t}\right) \omega_{f}  \tag{15}\\
\omega_{f} & =\left(\frac{I_{t}}{I_{t}+I_{d}}\right) \omega_{i} \tag{16}
\end{align*}
$$

The ratio of final to initial kinetic energy is then readily found:

$$
\begin{equation*}
\frac{K_{f}}{K_{i}}=\frac{\frac{1}{2}\left(I_{t}+I_{d}\right) \omega_{f}^{2}}{\frac{1}{2} I_{t} \omega_{i}^{2}}=\frac{I_{t}}{I_{t}+I_{d}} \tag{17}
\end{equation*}
$$

The fraction lost is then

$$
\begin{equation*}
\frac{K_{f}-K_{i}}{K_{i}}=\frac{K_{f}}{K_{i}}-1=\frac{-I_{d}}{I_{t}+I_{d}}=\frac{-0.54}{1+0.54} \approx-0.35 \tag{18}
\end{equation*}
$$

Approximately $35 \%$ of the initial kinetic energy is lost.
5. A string is wrapped around a pulley with a radius of 2.0 cm and no appreciable friction in its axle. The pulley is initially not turning. A constant force of 50 N is applied to the string, which does not slip, causing the pulley to rotate and the string to unwind. If the string unwinds 1.2 m in 4.9 s , what is the moment of inertia of the pulley?

Solution: The distance the string travels $\Delta x$ in a time $t$ implies an acceleration $a$ :

$$
\begin{equation*}
\Delta x=\frac{1}{2} a t^{2} \tag{19}
\end{equation*}
$$

Since the string does not slip, the rotational acceleration of the pulley must match the linear acceleration of the string divided by the radius of the pulley, $\alpha=a / R$. A torque balance relates the acceleration to the force present:

Name \& ID

$$
\begin{align*}
\sum \tau & =R F=I \alpha=I \frac{a}{R} \quad \Longrightarrow \quad a=\frac{R^{2} F}{I}  \tag{20}\\
\Delta x & =\frac{1}{2} a t^{2}=\frac{R^{2} F}{2 I} t^{2}  \tag{21}\\
I & =\frac{R^{2} F t^{2}}{2 \Delta x} \approx 0.20 \mathrm{~kg} \cdot \mathrm{~m}^{2} \tag{22}
\end{align*}
$$

## Formula sheet

basics

$$
g=\left|\overrightarrow{\mathbf{a}}_{\text {free fall }}\right|=9.81 \mathrm{~m} / \mathrm{s}^{2} \quad \text { near earth's surface }
$$

sphere $\quad V=\frac{4}{3} \pi r^{3}$

$$
a x^{2}+b x^{2}+c=0 \Longrightarrow x=\frac{-b \pm \sqrt{b^{2}-4 a c}}{2 a}
$$

## 1D \& 2D motion

$$
\begin{aligned}
\text { speed } & =v=|\overrightarrow{\mathbf{v}}| \quad \overrightarrow{\mathbf{v}}_{a v} \equiv \frac{\Delta \overrightarrow{\mathbf{r}}}{\Delta t} \quad \overrightarrow{\mathbf{v}}=\lim _{\Delta t \rightarrow 0} \frac{\Delta \overrightarrow{\mathbf{r}}}{\Delta t} \equiv \frac{d \overrightarrow{\mathbf{r}}}{d t} \\
a_{x} & =\lim _{\Delta t \rightarrow 0} \frac{\Delta v_{x}}{\Delta t} \equiv \frac{d v_{x}}{d t}=\frac{d}{d t}\left(\frac{d x}{d t}\right)=\frac{d^{2} x}{d t^{2}} \\
v_{x}(t) & =v_{x, i}+a_{x} t \\
x(t) & =x_{i}+v_{x, i} t+\frac{1}{2} a_{x} t^{2} \\
v_{x, f}^{2} & =v_{x, i}^{2}+2 a_{x} \Delta x
\end{aligned}
$$

$\downarrow$ launched from origin, level ground

$$
\begin{aligned}
y(x) & =\left(\tan \theta_{o}\right) x-\frac{g x^{2}}{2 v_{o}^{2} \cos ^{2} \theta_{o}} \\
\max \text { height } & =H=\frac{v_{i}^{2} \sin ^{2} \theta_{i}}{2 g} \\
\text { Range } & =R=\frac{v_{i}^{2} \sin 2 \theta_{i}}{g}
\end{aligned}
$$

## interactions

$$
\begin{aligned}
\Delta U^{G} & =m g \Delta x \quad \frac{a_{1 x}}{a_{2 x}}=-\frac{m_{2}}{m_{1}} \\
E_{\text {mech }} & =K+U \quad K=\frac{1}{2} m v^{2} \\
\Delta E_{\text {mech }} & =\Delta K+\Delta U=0 \quad \text { non-dissipative, closed } \\
\Delta E & =W \\
\Delta U_{\text {spring }} & =\frac{1}{2} k\left(x-x_{o}\right)^{2} \\
P & =\frac{d E}{d t} \text { general } \quad P=F_{\mathrm{ext}, \mathrm{x}} v_{x} \quad \text { 1D const force }
\end{aligned}
$$

## Rotation: we use radians

$$
\begin{aligned}
s & =\theta r \quad \leftarrow \text { arclength } \\
\omega & =\frac{d \theta}{d t}=\frac{v_{t}}{r} \quad \alpha=\frac{d \omega}{d t} \\
a_{t} & =\alpha r \quad \text { tangential } \quad a_{r}=-\frac{v^{2}}{r}=-\omega^{2} r \quad \text { radial } \\
v_{t} & =r \omega \quad v_{r}=0 \\
\Delta \theta & =\omega_{i} t+\frac{1}{2} \alpha t^{2} \quad \text { const } \alpha \\
\omega_{f} & =\omega_{i}+\alpha t \quad \text { const } \alpha \\
\Delta x & =r \theta \quad v=r \omega \quad a=r \alpha \quad \text { no slipping }
\end{aligned}
$$

$$
\begin{aligned}
I & =\sum_{i} m_{i} r_{i}^{2} \Rightarrow \int r^{2} d m=c m r^{2} \quad I=m r^{2} \quad \text { point particle } \\
I_{z} & =I_{c o m}+m d^{2} \quad \text { axis } z \text { parallel, dist } d \\
\overrightarrow{\mathbf{L}} & =\overrightarrow{\mathbf{r}} \times \overrightarrow{\mathbf{p}}=I \overrightarrow{\boldsymbol{\omega}} \quad L=I \omega=m v r_{\perp} \\
K & =\frac{1}{2} I \omega^{2}=L^{2} / 2 I \\
\Delta K & =\frac{1}{2} I \omega_{f}^{2}-\frac{1}{2} I \omega_{i}^{2}=W=\int \tau d \theta \\
P & =\frac{d W}{d t}=\tau \omega \\
\tau & =r F \sin \theta_{r F}=r_{\perp} F=r F_{\perp} \\
\tau_{n e t} & =\sum \vec{\tau}=I \overrightarrow{\boldsymbol{\alpha}}=\frac{d \mathbf{\mathbf { L }}}{d t} \\
K_{\mathrm{tot}} & =K_{c m}+K_{r o t}=\frac{1}{2} m v_{c m}^{2}+\frac{1}{2} I \omega^{2}
\end{aligned}
$$

$$
\begin{aligned}
& \text { work } \\
& \begin{aligned}
\Delta E_{\mathrm{mech}} & =\Delta K+\Delta U=W \quad \leftarrow \text { not closed } \quad \Delta U_{\text {spring }}=\frac{1}{2} k\left(x-x_{o}\right)^{2} \\
P & =\frac{d E}{d t} \quad P=F_{\text {ext }, \mathrm{x}} v_{x} \quad \text { one dimension } \\
W & =\left(\sum^{\overrightarrow{\mathbf{F}}}\right) \Delta x_{F} \quad \text { constant foce 1D } \\
W & =\sum_{n}\left(F_{\text {ext }, \mathrm{x}} \Delta x_{F n}\right) \quad \text { const nondiss., many particles, 1D } \\
W & =\int_{x_{i}}^{x_{f}} F_{x}(x) d x \quad \text { nondiss. force, 1D } \\
\left(F_{12}^{s}\right)_{\max } & =\mu_{s} F_{12}^{n} \quad \text { static } \quad F_{12}^{k}=\mu_{k} F_{12}^{n} \quad \text { kinetic } \\
W & =\overrightarrow{\mathbf{F}}^{2} \cdot \Delta \overrightarrow{\mathbf{r}}_{F} \quad \text { const non-diss force } \\
W & =\int_{f} \overrightarrow{\mathbf{F}}(\overrightarrow{\mathbf{r}}) \cdot d \overrightarrow{\mathbf{r}} \quad \text { variable nondiss force } \\
& \overrightarrow{\mathbf{r}}_{i}
\end{aligned}
\end{aligned}
$$

Moments of inertia of things of mass $M$

| Object | axis | dimension | $\mathbf{I}$ |
| :--- | :---: | :---: | :---: |
| solid sphere | central axis | radius $R$ | $\frac{2}{5} M R^{2}$ |
| hollow sphere | central axis | radius $R$ | $\frac{2}{3} M R^{2}$ |
| solid disc/cylinder | central axis | radius $R$ | $\frac{1}{2} M R^{2}$ |
| hoop | central axis | radius $R$ | $M R^{2}$ |
| point particle | pivot point | distance $R$ to pivot | $M R^{2}$ |
| rod | center | length $L$ | $\frac{1}{12} M L^{2}$ |
| rod | end | length $L$ | $\frac{1}{3} M L^{2}$ |
| solid regular octahedron | through vertices | side $a$ | $\frac{1}{10} m a^{2}$ |


| Derived unit | Symbol | equivalent to |
| :--- | :---: | :---: |
| newton | N | $\mathrm{kg} \cdot \mathrm{m} / \mathrm{s}^{2}$ |
| joule | J | $\mathrm{kg} \cdot \mathrm{m}^{2} / \mathrm{s}^{2}=\mathrm{N} \cdot \mathrm{m}$ |
| watt | W | $\mathrm{J} / \mathrm{s}=\mathrm{m}^{2} \cdot \mathrm{~kg} / \mathrm{s}^{3}$ |

