# University of Alabama <br> Department of Physics and Astronomy 

PH 301 / LeClair

## Exam 1 redux - solution

1. A particle of mass $m$ traveling along a horizontal line with constant velocity $v_{o}$ suddenly experiences a resistive force $f=-b v^{3}$. (a) Find an expression for $v(t)$. Hint: it is always the same method. (b) Does the object eventually halt, either at finite time or as $t \rightarrow \infty$ ? If so, under what condition, and if not, why?
bonus $+\mathbf{1}$ : What is the time constant of the $v(t)$ decay? That is, your result should have a term that looks like $t / \tau$, what is $\tau$ ?

Solution: The only force present is $f$, so we just need to set up the equation of motion, separate variables, and integrate.

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\begin{align*}
m \ddot{x} & =m \dot{v}=m \frac{d v}{d t}=f=-b v^{3}  \tag{1}\\
\frac{d v}{v^{3}} & =-\frac{b}{m} d t  \tag{2}\\
\int_{v_{o}}^{v} \frac{d v}{v^{3}} & =-\int_{0}^{t} \frac{b}{m} d t=-\frac{b}{m} t  \tag{3}\\
-\frac{b}{m} t & =-\left.\frac{1}{2 v^{2}}\right|_{v_{o}} ^{v}=-\frac{1}{2}\left(\frac{1}{v^{2}}-\frac{1}{v_{o}^{2}}\right)  \tag{4}\\
\frac{2 b t}{m} & =\frac{1}{v^{2}}-\frac{1}{v_{o}^{2}}  \tag{5}\\
\frac{1}{v^{2}} & =\frac{2 b t}{m}+\frac{1}{v_{o}^{2}}  \tag{6}\\
v & =\sqrt{\frac{1}{\frac{2 b t}{m}+\frac{1}{v_{o}^{2}}}} \tag{7}
\end{align*}
$$

One can see that $v \rightarrow 0$ as $t \rightarrow \infty$, so the object does halt eventually. We can also see that $v$ is a function of $2 b t / m$. Comparing to the usual form $t / \tau$, we identify the time constant for velocity decay as $\tau=m / 2 b$.
2. A rectangular block of height $h$ and area $A$ floats in water. If the density of the block is $\rho_{b}$ and the density of the water is $\rho_{w}$, find the frequency of small oscillations when the block bobs up and down on the surface of the water. Hints: The buoyant force is the weight of the displaced fluid. If the object is floating in equilibrium, what is the net force if it is displaced by some $\delta x$ ?
bonus, $+\mathbf{2}$ : show that the frequency of oscillation is equivalent to that of a pendulum of length $d_{o}$, where $d_{o}$ is the submerged depth of the block in equilibrium.

Solution: First let's find the equilibrium condition. In equilibrium, the buoyant force (weight of displaced water) equals the block's weight. Let $x$ be length of the block that is below the surface of the water. Then

$$
\begin{align*}
F_{B} & =m g  \tag{8}\\
\rho_{w} A x g & =\rho_{b} A h g  \tag{9}\\
x_{\mathrm{eq}} & =\frac{\rho_{b}}{\rho_{w}} h \tag{10}
\end{align*}
$$

Now, if we push the block down an additional distance $\delta x$ from equilibrium, there will be a net force due to the new volume of water displaced $(A \delta x)$ :

$$
\begin{align*}
\delta F & =-\rho_{w} A g \delta x=m a=\rho_{b} A h a  \tag{11}\\
a & =-\frac{\rho_{w} g}{\rho_{b} h} \delta x=-\omega^{2} \delta x \quad \Longrightarrow \quad \omega=\sqrt{\frac{\rho_{w} g}{\rho_{b} h}} \tag{12}
\end{align*}
$$

We end up with simple harmonic motion about the equilibrium point. Noting that $x_{\text {eq }}=\frac{\rho_{b}}{\rho_{w}} h$, we can write $\omega=\sqrt{g / x_{\text {eq }}}$, meaning the block behaves like a pendulum whose length is equal to the equilibrium depth.
3. (a) Write down the total energy $E$ for a mass on a spring in 1 D in terms of $\dot{x}$ and $x$. (b) Show, as is the case for any conservative 1D system, that you can obtain the equation of motion for the coordinate $x$ by differentiating the equation $E=$ const. It should be the familiar Newton's 2nd law result.
bonus $+\mathbf{1}$ : Suppose the energy were not constant, but you had it on good authority that energy was lost to the environment at a rate proportional to velocity squared, i.e,. $\dot{E} \propto-\dot{x}^{2}$. With the result of (b), show that this is in fact just a damped harmonic oscillator. Hint: consider a constant of proportionality $2 \beta$ to produce a familiar-looking result.

Solution: We know the energy well enough. Remember the chain rule and $\dot{E}$ is found readily.

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\begin{align*}
E & =\frac{1}{2} m \dot{x}^{2}+\frac{1}{2} k x^{2}  \tag{13}\\
\dot{E} & =m \dot{x} \ddot{x}+k x \dot{x}=0  \tag{14}\\
m \ddot{x}+k x & =0 \tag{15}
\end{align*}
$$

Equation of motion recovered. If $\dot{E}=-2 \beta \dot{x}^{2}$,

$$
\begin{align*}
\dot{E} & =m \dot{x} \ddot{x}+k x \dot{x}=-2 \beta \dot{x}^{2}  \tag{16}\\
0 & =m \dot{x} \ddot{x}+2 \beta \dot{x}^{2}+k x \dot{x}  \tag{17}\\
0 & =m \ddot{x}+2 \beta \dot{x}+k x \tag{18}
\end{align*}
$$

This is the equation for a damped harmonic oscillator we are familiar with.

